

Twisted string bordism in 7 dimensions with applications to anomaly cancellation

Natalia María Pacheco-Tallaj

nataliap@mit.edu

based on joint work with I. Basile, C. Krulewski, G. Leone

Chicago Topology Seminar

Outline

- 1 Anomalies in quantum field theory
 - Quantization, symmetries, and anomalies
 - Local anomalies and twisted tangential structures
 - Anomalies as invertible bulk theories
- 2 Invertible field theories and bordism
 - Functorial field theory
 - IFTs and stable homotopy theory
 - Our mathematical setup
- 3 Twists of $BU(1)$ -string bordism in dimension 7
- 4 An illustrative example: a bordism invariant for $n = 1$
- 5 Manifold generators

Anomalies in quantum field theory

What is a field theory?

A **field theory** is the data of

- A *spacetime manifold* X which is d dimensional
- A moduli $\mathcal{F}(X)$ of objects over X , called *fields*, e.g.
 - connections on a principal G -bundle $P \rightarrow X$ (gauge fields)
 - sections of a vector bundle $E \rightarrow X$ (matter fields)
 - maps to a target manifold $\text{Maps}(X, Y)$ (sigma models)
 - metrics on X (gravity)
- An *action functional* $S : \mathcal{F}(X) \rightarrow \mathbb{C}$

What is a classical field theory?

A **field theory** is the data of

- A *spacetime manifold* X which is d dimensional
- A moduli $\mathcal{F}(X)$ of objects over X , called *fields*, e.g.
 - connections on a principal G -bundle $P \rightarrow X$ (gauge fields)
 - sections of a vector bundle $E \rightarrow X$ (matter fields)
 - maps to a target manifold $\text{Maps}(X, Y)$ (sigma models)
 - metrics on X (gravity)
- An *action functional* $S : \mathcal{F}(X) \rightarrow \mathbb{C}$

A **classical field theory** finds physical fields ψ by solving a principle of stationary action $\delta S / \delta \psi = 0$.

What is a quantum field theory?

A **field theory** is the data of

- A *spacetime manifold* X which is d dimensional
- A moduli $\mathcal{F}(X)$ of objects over X , called *fields*, e.g.
 - connections on a principal G -bundle $P \rightarrow X$ (gauge fields)
 - sections of a vector bundle $E \rightarrow X$ (matter fields)
 - maps to a target manifold $\text{Maps}(X, Y)$ (sigma models)
 - metrics on X (gravity)
- An *action functional* $S : \mathcal{F}(X) \rightarrow \mathbb{C}$

A **quantum field theory** instead defines a probability measure on $\mathcal{F}(X)$ weighted by $\exp(-S)$ (assuming X has Euclidean signature) and computes probability amplitudes and correlation functions of observable physical quantities in this measure.

Partition functions and path integrals

- QFT: probability distribution over space of fields weighted by $\exp(-S)$
- Path integral/partition function: integrate over **dynamical fields** ψ and consider result as a function of the **background fields** A

$$Z_X[A] = \int_{\psi \in \mathcal{F}^{\text{dyn}}(X)} D\psi e^{-S[\psi, A]} : \mathcal{F}^{\text{bg}}(X) \rightarrow \mathbb{C}$$

Partition functions and path integrals

- QFT: probability distribution over space of fields weighted by $\exp(-S)$
- Path integral/partition function: integrate over **dynamical fields** ψ and consider result as a function of the **background fields** A

$$Z_X[A] = \int_{\psi \in \mathcal{F}^{\text{dyn}}(X)} D\psi e^{-S[\psi, A]} : \mathcal{F}^{\text{bg}}(X) \rightarrow \mathbb{C}$$

(schematically $\mathcal{F}^{\text{dyn}}(X) \rightarrow \mathcal{F}(X) \rightarrow \mathcal{F}^{\text{bg}}(X)$)

Partition functions and path integrals

- QFT: probability distribution over space of fields weighted by $\exp(-S)$
- Path integral/partition function: integrate over **dynamical fields** ψ and consider result as a function of the **background fields** A

$$Z_X[A] = \int_{\psi \in \mathcal{F}^{\text{dyn}}(X)} D\psi e^{-S[\psi, A]} : \mathcal{F}^{\text{bg}}(X) \rightarrow \mathbb{C}$$

$$(\text{schematically } \underbrace{\mathcal{F}^{\text{dyn}}(X)}_{\text{fermions, matter}} \rightarrow \mathcal{F}(X) \rightarrow \underbrace{\mathcal{F}^{\text{bg}}(X)}_{\text{background gauge fields}})$$

What is a quantum anomaly?

“A failure of $Z[A]$ to be defined globally and gauge invariantly.”

Example

- X : $4d$ -dimensional spin manifold with a principal G -bundle $P \rightarrow X$ and associated bundle $E_P := P \times_G V$

What is a quantum anomaly?

“A failure of $Z[A]$ to be defined globally and gauge invariantly.”

Example

- X : $4d$ -dimensional spin manifold with a principal G -bundle $P \rightarrow X$ and associated bundle $E_P := P \times_G V$
- $\mathcal{A} := \text{Conn}(P)$ - space of connections, $\mathcal{F}_{\text{fer}} := \Gamma(X, S^+ \otimes E_P)$ - space of spinors, so $\mathcal{F}(X) = (\mathcal{A} \times \mathcal{F}_{\text{fer}})/\mathcal{G}$

What is a quantum anomaly?

“A failure of $Z[A]$ to be defined globally and gauge invariantly.”

Example

- X : $4d$ -dimensional spin manifold with a principal G -bundle $P \rightarrow X$ and associated bundle $E_P := P \times_G V$
- $\mathcal{A} := \text{Conn}(P)$ - space of connections, $\mathcal{F}_{\text{fer}} := \Gamma(X, S^+ \otimes E_P)$ - space of spinors, so $\mathcal{F}(X) = (\mathcal{A} \times \mathcal{F}_{\text{fer}})/\mathcal{G}$
- Action functional $S[A, \psi] := \int_X \psi^\dagger D_A \psi$

What is a quantum anomaly?

“A failure of $Z[A]$ to be defined globally and gauge invariantly.”

Example

- X : $4d$ -dimensional spin manifold with a principal G -bundle $P \rightarrow X$ and associated bundle $E_P := P \times_G V$
- $\mathcal{A} := \text{Conn}(P)$ - space of connections, $\mathcal{F}_{\text{fer}} := \Gamma(X, S^+ \otimes E_P)$ - space of spinors, so $\mathcal{F}(X) = (\mathcal{A} \times \mathcal{F}_{\text{fer}})/\mathcal{G}$
- Action functional $S[A, \psi] := \int_X \psi^\dagger D_A \psi$
- Integrate out ψ : $Z[A] := \int_{\psi \in \Gamma(X, S^+ \otimes E_A)} e^{-\langle \psi, D_A \psi \rangle}$ is a section of the determinant line bundle $\mathcal{L}_{\text{anom}}$ over the moduli stack $[\mathcal{A}/\mathcal{G}]$, associated to the family of chiral Dirac operators $\{D_A\}$ on X .

What is a quantum anomaly?

“A failure of $Z[A]$ to be defined globally and gauge invariantly.”

Example

- X : $4d$ -dimensional spin manifold with a principal G -bundle $P \rightarrow X$ and associated bundle $E_P := P \times_G V$
- $\mathcal{A} := \text{Conn}(P)$ - space of connections, $\mathcal{F}_{\text{fer}} := \Gamma(X, S^+ \otimes E_P)$ - space of spinors, so $\mathcal{F}(X) = (\mathcal{A} \times \mathcal{F}_{\text{fer}})/\mathcal{G}$
- Action functional $S[A, \psi] := \int_X \psi^\dagger D_A \psi$
- Integrate out ψ : $Z[A] := \int_{\psi \in \Gamma(X, S^+ \otimes E_A)} e^{-\langle \psi, D_A \psi \rangle}$ is a section of the determinant line bundle $\mathcal{L}_{\text{anom}}$ over the moduli stack $[\mathcal{A}/\mathcal{G}]$, associated to the family of chiral Dirac operators $\{D_A\}$ on X .
- By the family index theorem, we have a class $P_{d+2} := [\hat{A}(X) \text{ch}(F)]_{(d+2)}$ on $X \times [\mathcal{A}/\mathcal{G}]$ such that $\int_X P_{d+2} \in \Omega^2(\mathcal{A}/\mathcal{G})$ is the curvature form of $\mathcal{L}_{\text{anom}}$.

Local anomalies are encoded in anomaly polynomials

$$\mathcal{F}^{\text{dyn}}(X) \rightarrow \mathcal{F}(X) \rightarrow \mathcal{F}^{\text{bg}}(X)$$

$$Z_X[A] = \int_{\psi \in \mathcal{F}^{\text{dyn}}(X)} D\psi e^{-S[\psi, A]} \in \Gamma(\mathcal{F}^{\text{bg}}(X), \mathcal{L}_{\text{anom}})$$

- Families index theorem $\implies$ curvature of $\mathcal{L}_{\text{anom}}$ is encoded by a degree $d + 2$ **anomaly polynomial**
 - an index quantity on X
 - characteristic forms of the gauge fields

Local anomalies are encoded in anomaly polynomials

$$\mathcal{F}^{\text{dyn}}(X) \rightarrow \mathcal{F}(X) \rightarrow \mathcal{F}^{\text{bg}}(X)$$

$$Z_X[A] = \int_{\psi \in \mathcal{F}^{\text{dyn}}(X)} D\psi e^{-S[\psi, A]} \in \Gamma(\mathcal{F}^{\text{bg}}(X), \mathcal{L}_{\text{anom}})$$

- Families index theorem $\implies$ curvature of $\mathcal{L}_{\text{anom}}$ is encoded by a degree $d+2$ **anomaly polynomial**
 - an index quantity on X
 - characteristic forms of the gauge fields
- (Green-Schwarz mechanism) If $P_{d+2} = X_4 \wedge X_{d-2}$, can construct a “counterterm” to cancel the anomaly

Anomaly inflow

- (Green-Schwarz mechanism) If $P_{d+2} = X_4 \wedge X_{d-2}$, can construct a “counterterm” α_{anom} to cancel the anomaly

$$Z_X[A]e^{-2\pi i\alpha_{\text{anom}}}$$

- (Anomaly inflow) This phase comes from the partition function of a field theory in dimension $d+1$: there is some M with $\partial M = X$ and such that

$$Z_X[A]e^{-2\pi i(\text{Id}_X(M) + \int_M H \wedge X_{d-2})}$$
 is gauge invariant

only well defined if X_4 is cohomologically trivial on M .

- We call this $d+1$ -dimensional theory the **anomaly theory** of Z

Our setup: 6d supergravity theories

- We consider 6d $\mathcal{N} = (1, 0)$ supergravity theories with type A-D-E or abelian gauge groups

Our setup: 6d supergravity theories

- We consider 6d $\mathcal{N} = (1, 0)$ supergravity theories with $U(1)$ gauge group, meaning our 6-dimensional manifolds X are equipped with $f : X \rightarrow BU(1)$ inducing a complex line bundle $\mathcal{L} := f^*(\mathcal{O}(1))$.

Our setup: 6d supergravity theories

- We consider 6d $\mathcal{N} = (1, 0)$ supergravity theories with $U(1)$ gauge group, meaning our 6-dimensional manifolds X are equipped with $f : X \rightarrow BU(1)$ inducing a complex line bundle $\mathcal{L} := f^*(\mathcal{O}(1))$.
- Anomaly polynomial is

$$(\frac{1}{2}p_1(X) - nc_1(\mathcal{L})^2)c_1(\mathcal{L})^2$$

Our setup: 6d supergravity theories

- We consider 6d $\mathcal{N} = (1, 0)$ supergravity theories with $U(1)$ gauge group, meaning our 6-dimensional manifolds X are equipped with $f : X \rightarrow BU(1)$ inducing a complex line bundle $\mathcal{L} := f^*(\mathcal{O}(1))$.
- Anomaly polynomial is

$$\left(\frac{1}{2}p_1(X) - nc_1(\mathcal{L})^2\right)c_1(\mathcal{L})^2$$

- Upshot: all the information of the anomalies is encoded in a $7d$ anomaly theory defined on 7-manifolds with $\frac{1}{2}p_1(TM) = nc_1(\mathcal{L})^2$

String structures

7-manifolds with $\frac{1}{2}p_1(TM) = nc_1(\mathcal{L})^2$

Definition (String structure)

Let M be a spin manifold, then $p_1(M)$ is even. A *string structure* on M is a trivialization of $\frac{1}{2}p_1(M)$.

$$\begin{array}{ccccc} & & BString & & \\ & \nearrow & \downarrow & & \\ M & \xrightarrow{TM} & BSpin & \xrightarrow{\frac{1}{2}p_1} & K(\mathbb{Z}, 4) \end{array}$$

Twisted string structures

7-manifolds with $\frac{1}{2}p_1(TM) = nc_1(\mathcal{L})^2$

Definition (String structure)

Let M be a spin manifold, then $p_1(M)$ is even. A *string structure* on M is a trivialization of $\frac{1}{2}p_1(M)$.

Definition (nc_1^2 -twisted string structure)

Let M be a spin manifold equipped with a map $f : M \rightarrow BU(1)$. An nc_1^2 -twisted string structure is a trivialization of $\frac{1}{2}p_1(M) - nc_1(\mathcal{L})^2$.

$$\begin{array}{ccc}
 M & \xrightarrow{f} & BU(1) \\
 \tau M \downarrow & & \downarrow nc_1(\mathcal{O}(1))^2 \\
 BSpin & \xrightarrow{\frac{1}{2}p_1} & K(\mathbb{Z}, 4)
 \end{array}$$

Twisted string structures

7-manifolds with $\frac{1}{2}p_1(TM) = nc_1(\mathcal{L})^2$

Definition (nc_1^2 -twisted string structure)

Let M be a spin manifold equipped with a map $f : M \rightarrow BU(1)$.
 An nc_1^2 -twisted string structure is a trivialization of
 $\frac{1}{2}p_1(M) - nc_1(\mathcal{L})^2$.

Twisted string structures

7-manifolds with $\frac{1}{2}p_1(TM) = nc_1(\mathcal{L})^2$

Definition (nc_1^2 -twisted string structure)

Let M be a spin manifold equipped with a map $f : M \rightarrow BU(1)$.
An nc_1^2 -twisted string structure is a trivialization of $\frac{1}{2}p_1(M) - nc_1(\mathcal{L})^2$.

Remark

Consider the (virtual) bundle $T = \mathcal{O}(1) + \mathcal{O}(-1) - 2\mathbb{C}$ on $BU(1)$. We also call this a $(BU(1), -nT)$ -twisted string structure, because given a manifold M with a map $f : M \rightarrow BU(1)$, it is the data of a string structure on $TM + f^(-nT)$.^a*

$$^a \frac{1}{2}p_1(TM - f^*nT) = \frac{1}{2}p_1(TM) - nc_1(f^*\mathcal{O}(1))^2 = 0$$

Summary

- QFT: Try to integrate $D\psi e^{-S[A,\psi]}$ over dynamical fields, result is section of $\mathcal{L}_{\text{anom}} \rightarrow \mathcal{F}^{\text{bg}}(X)$

Summary

- QFT: Try to integrate $D\psi e^{-S[A,\psi]}$ over dynamical fields, result is section of $\mathcal{L}_{\text{anom}} \rightarrow \mathcal{F}^{\text{bg}}(X)$
- Local anomaly (curvature of $\mathcal{L}_{\text{anom}}$) encoded in $P_{d+2} = X_4 \wedge X_{d-4}$

Summary

- QFT: Try to integrate $D\psi e^{-S[A,\psi]}$ over dynamical fields, result is section of $\mathcal{L}_{\text{anom}} \rightarrow \mathcal{F}^{\text{bg}}(X)$
- Local anomaly (curvature of $\mathcal{L}_{\text{anom}}$) encoded in $P_{d+2} = X_4 \wedge X_{d-4}$
- $Z_X[A]$ not gauge invariant but $Z_X[A]e^{-2\pi i \alpha_M[\tilde{A}]}$ is, with $\partial M = X$

Summary

- QFT: Try to integrate $D\psi e^{-S[A,\psi]}$ over dynamical fields, result is section of $\mathcal{L}_{\text{anom}} \rightarrow \mathcal{F}^{\text{bg}}(X)$
- Local anomaly (curvature of $\mathcal{L}_{\text{anom}}$) encoded in $P_{d+2} = X_4 \wedge X_{d-4}$
- $Z_X[A]$ not gauge invariant but $Z_X[A]e^{-2\pi i \alpha_M[\tilde{A}]}$ is, with $\partial M = X$
- α_M is a well-defined field theory in $d+1$ dimensions if M has $X_4 = 0$, which is a twisted tangential structure

Summary

- QFT: Try to integrate $D\psi e^{-S[A,\psi]}$ over dynamical fields, result is section of $\mathcal{L}_{\text{anom}} \rightarrow \mathcal{F}^{\text{bg}}(X)$
- Local anomaly (curvature of $\mathcal{L}_{\text{anom}}$) encoded in $P_{d+2} = X_4 \wedge X_{d-4}$
- $Z_X[A]$ not gauge invariant but $Z_X[A]e^{-2\pi i \alpha_M[\tilde{A}]}$ is, with $\partial M = X$
- α_M is a well-defined field theory in $d+1$ dimensions if M has $X_4 = 0$, which is a twisted tangential structure
- We call it the *anomaly theory*, and it is an invertible topological field theory defined on X_4 manifolds

Summary

- QFT: Try to integrate $D\psi e^{-S[A,\psi]}$ over dynamical fields, result is section of $\mathcal{L}_{\text{anom}} \rightarrow \mathcal{F}^{\text{bg}}(X)$
- Local anomaly (curvature of $\mathcal{L}_{\text{anom}}$) encoded in $P_{d+2} = (\frac{1}{2}p_1(X) - nc_1(\mathcal{L})^2)c_1(\mathcal{L})^2$
- $Z_X[A]$ not gauge invariant but $Z_X[A]e^{-2\pi i\alpha_M[\tilde{A}]}$ is, with $\partial M = X$
- α_M is a well-defined field theory in $d + 1$ dimensions if M is nc_1^2 -twisted string
- We call it the *anomaly theory*, and it is an invertible topological field theory defined on nc_1^2 -twisted string manifolds

Invertible field theories and bordism

The functorial formalism and its interpretation

The data of a field theory in $d + 1$ dimensions can often be organized mathematically as follows:

Definition (Functorial field theory)

A $d + 1$ -dimensional functorial field theory (on nc_1^2 -twisted String manifolds) is a symmetric monoidal functor

$$Z : (\text{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-String}}, \sqcup) \rightarrow (\text{sVect}_{\mathbb{C}}, \otimes)$$

The functorial formalism and its interpretation

Definition (Functorial field theory)

A $d + 1$ -dimensional functorial field theory (on nc_1^2 -twisted String manifolds) is a symmetric monoidal functor

$$Z : (\text{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-String}}, \sqcup) \rightarrow (s\text{Vect}_{\mathbb{C}}, \otimes)$$

- $Z(X^d)$ is interpreted as a state space of the theory
- For a closed $d + 1$ manifold M , interpreted as a bordism $\emptyset \rightsquigarrow \emptyset$, the morphism $Z(M) : \mathbb{C} \rightarrow \mathbb{C}$ in $s\text{Vect}_{\mathbb{C}}$ is determined by a choice of complex number $\eta \in \mathbb{C}$ interpreted as the value of the partition function on M

Invertibility

There is a tensor product structure on the set of all functorial field theories given by point-wise tensor product in $(\mathbf{sVect}_{\mathbb{C}}, \otimes)$:

$$(Z_1 \otimes Z_2)(X^d) = Z_1(X) \otimes Z_2(X)$$

Invertibility

There is a tensor product structure on the set of all functorial field theories given by point-wise tensor product in $(s\text{Vect}_{\mathbb{C}}, \otimes)$:

$$(Z_1 \otimes Z_2)(X^d) = Z_1(X) \otimes Z_2(X)$$

Definition (Invertible field theory)

A functorial field theory $Z : \text{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-String}} \rightarrow s\text{Vect}_{\mathbb{C}}$ is said to be *invertible* if it is tensor-invertible under the above tensor product.

Invertible field theories

Invertible field theories (IFTs) factor through the maximal Picard groupoid ¹ of the target $s\text{Line}_{\mathbb{C}} \subset s\text{Vect}_{\mathbb{C}}$

$$\begin{array}{ccc}
 (\text{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup) & \xrightarrow{Z} & (s\text{Vect}_{\mathbb{C}}, \otimes) \\
 & \searrow \text{dashed} & \uparrow \\
 & & (s\text{Line}_{\mathbb{C}}, \otimes)
 \end{array}$$

¹A Picard groupoid is a fully invertible symmetric monoidal category

Invertible field theories

Invertible field theories (IFTs) factor through the groupoid completion of the source.

$$\begin{array}{ccc}
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup) & \xrightarrow{Z} & (s\mathrm{Vect}_{\mathbb{C}}, \otimes) \\
 \downarrow & \searrow \text{dashed} & \uparrow \\
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup)^{\mathrm{gp}} & \dashrightarrow & (s\mathrm{Line}_{\mathbb{C}}, \otimes)
 \end{array}$$

Classifying invertible field theories

IFTs are classified by maps of spectra.

$$\begin{array}{ccc}
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup) & \xrightarrow{Z} & (s\mathrm{Vect}_{\mathbb{C}}, \otimes) \\
 \downarrow & \searrow \text{dashed} & \uparrow \\
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup)^{\mathrm{gp}} & \dashrightarrow & (s\mathrm{Line}_{\mathbb{C}}, \otimes)
 \end{array}$$

bordism spectrum $\longrightarrow$ character dual of the sphere

Classifying invertible field theories

IFTs are classified by maps of spectra.

$$\begin{array}{ccc}
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup) & \xrightarrow{Z} & (s\mathrm{Vect}_{\mathbb{C}}, \otimes) \\
 \downarrow & \searrow \text{dashed} & \uparrow \\
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup)^{\mathrm{gp}} & \dashrightarrow & (s\mathrm{Line}_{\mathbb{C}}, \otimes)
 \end{array}$$

$$\begin{array}{ccc}
 \text{bordism} & \longrightarrow & \text{character} \\
 \text{spectrum} & & \text{dual of the} \\
 & & \text{sphere}
 \end{array}$$

$$MString \wedge BU(1)^{-nT} \longrightarrow I\mathbb{C}^{\times}$$

Classifying invertible field theories

IFTs are classified by maps of spectra.

$$\begin{array}{ccc}
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup) & \xrightarrow{Z} & (s\mathrm{Vect}_{\mathbb{C}}, \otimes) \\
 \downarrow & \searrow \text{dashed} & \uparrow \\
 (\mathrm{Bord}_{\langle d, d+1 \rangle}^{nc_1^2\text{-twisted}}, \sqcup)^{\mathrm{gp}} & \dashrightarrow & (s\mathrm{Line}_{\mathbb{C}}, \otimes)
 \end{array}$$

$$\begin{array}{ccc}
 \text{bordism} & \xrightarrow{\quad} & \text{character} \\
 \text{spectrum} & & \text{dual of the} \\
 & & \text{sphere}
 \end{array}$$

$$M\mathrm{String} \wedge BU(1)^{-nT} \longrightarrow I\mathbb{C}^{\times}$$

Classification theorems: Freed-Hopkins, Freed-Hopkins-Teleman, Grady (deformation classes)

Summary of the physics $\rightsquigarrow$ topology conversion

- When we study anomaly cancellation of a d -dimensional theory, we:

Summary of the physics $\rightsquigarrow$ topology conversion

- When we study anomaly cancellation of a d -dimensional theory, we:
 - ① Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by
$$P_{d+2} = X_4 \wedge X_{d-2}$$

Summary of the physics $\rightsquigarrow$ topology conversion

- When we study anomaly cancellation of a d -dimensional theory, we:
 - 1 Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by $P_{d+2} = X_4 \wedge X_{d-2}$
 - 2 Represent the total anomaly by $(d+1)$ -dimensional IFT in terms of X_{d-2} with the twisted tangential structure determined by X_4

Summary of the physics $\rightsquigarrow$ topology conversion

- When we study anomaly cancellation of a d -dimensional theory, we:
 - ① Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by $P_{d+2} = X_4 \wedge X_{d-2}$
 - ② Represent the total anomaly by $(d+1)$ -dimensional IFT in terms of X_{d-2} with the twisted tangential structure determined by X_4
 - ③ IFTs are determined by the map they define out of a bordism group of manifolds with X_4 structure

Summary of the physics $\rightsquigarrow$ topology conversion

- When we study anomaly cancellation of a d -dimensional theory, we:
 - 1 Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by $P_{6+2} = (\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)) \wedge c_1(\mathcal{L})^2$
 - 2 Represent the total anomaly by $(d+1)$ -dimensional IFT in terms of $c_1(\mathcal{L})^2$ with the twisted tangential structure determined by $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)$
 - 3 IFTs are determined by the map they define out of a bordism group of manifolds with $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2) = 0$ structure

Goals

- When we study anomaly cancellation of a d -dimensional theory, we:
 - ① Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by $P_{6+2} = (\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)) \wedge c_1(\mathcal{L})^2$
 - ② Represent the total anomaly by $(d+1)$ -dimensional IFT in terms of $c_1(\mathcal{L})^2$ with the twisted tangential structure determined by $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)$
 - ③ IFTs are determined by the map they define out of a bordism group of manifolds with $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2) = 0$ structure
- Goal: prove anomaly cancellation in 6d supergravity with $U(1)$ gauge group

Goals

- When we study anomaly cancellation of a d -dimensional theory, we:
 - ① Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by $P_{6+2} = (\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)) \wedge c_1(\mathcal{L})^2$
 - ② Represent the total anomaly by $(d+1)$ -dimensional IFT in terms of $c_1(\mathcal{L})^2$ with the twisted tangential structure determined by $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)$
 - ③ IFTs are determined by the map they define out of a bordism group of manifolds with $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2) = 0$ structure
- Goal: prove anomaly cancellation in 6d supergravity with $U(1)$ gauge group
 - ① Compute 7-dimensional twisted string bordism groups.

Goals

- When we study anomaly cancellation of a d -dimensional theory, we:
 - ① Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by $P_{6+2} = (\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)) \wedge c_1(\mathcal{L})^2$
 - ② Represent the total anomaly by $(d+1)$ -dimensional IFT in terms of $c_1(\mathcal{L})^2$ with the twisted tangential structure determined by $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)$
 - ③ IFTs are determined by the map they define out of a bordism group of manifolds with $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2) = 0$ structure
- Goal: prove anomaly cancellation in 6d supergravity with $U(1)$ gauge group
 - ① Compute 7-dimensional twisted string bordism groups. They are nonzero!

Goals

- When we study anomaly cancellation of a d -dimensional theory, we:
 - ① Compute the local anomaly: curvature of $\mathcal{L}_{\text{anom}}$, detected by $P_{6+2} = (\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)) \wedge c_1(\mathcal{L})^2$
 - ② Represent the total anomaly by $(d+1)$ -dimensional IFT in terms of $c_1(\mathcal{L})^2$ with the twisted tangential structure determined by $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2)$
 - ③ IFTs are determined by the map they define out of a bordism group of manifolds with $\frac{1}{2}p_1(X) - nc_1(\mathcal{L}^2) = 0$ structure
- Goal: prove anomaly cancellation in 6d supergravity with $U(1)$ gauge group
 - ① Compute 7-dimensional twisted string bordism groups. They are nonzero!
 - ② Construct 7-dimensional manifold generators to compute α_{anom} on

Twists of $BU(1)$ -string bordism in dimension 7

Challenges computing $\Omega_7^{\text{String}-nc_1^2}$

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.

Challenges computing $\Omega_7^{\text{String}-nc_1^2}$

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$

Challenges computing $\Omega_7^{\text{String}-nc_1^2}$

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups.

Challenges computing $\Omega_7^{\text{String}-nc_1^2}$

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups. Smith homomorphisms provide limited success for some $-nT$ twists but we want something that works for all $-nT$ twists.

Challenges computing $\Omega_7^{\text{String}-nc_1^2}$

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups. Smith homomorphisms provide limited success for some $-nT$ twists but we want something that works for all $-nT$ twists.
- Generators?
 - $K3 \times S^3$, $\mathbb{CP}^2 \times S^3$ are not twisted string,

Challenges computing $\Omega_7^{\text{String}-nc_1^2}$

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups. Smith homomorphisms provide limited success for some $-nT$ twists but we want something that works for all $-nT$ twists.
- Generators?
 - $K3 \times S^3$, $\mathbb{CP}^2 \times S^3$ are not twisted string,
 - odd dimension means no hope for complex algebraic representatives.

Challenges computing $\Omega_7^{\text{String}-nc_1^2}$

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups. Smith homomorphisms provide limited success for some $-nT$ twists but we want something that works for all $-nT$ twists.
- Generators?
 - $K3 \times S^3$, $\mathbb{CP}^2 \times S^3$ are not twisted string,
 - odd dimension means no hope for complex algebraic representatives.
 - Lens spaces? Can determine twisted string structure, but hard to detect if generators.

Solutions

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups. Smith homomorphisms provide limited success for some $-nT$ twists but we want something that works for all $-nT$ twists.
- Generators?
 - $K3 \times S^3$, $\mathbb{CP}^2 \times S^3$ are not twisted string,
 - odd dimension means no hope for complex algebraic representatives.
 - Lens spaces? Can determine twisted string structure, but hard to detect if generators.

Solutions: characterize homotopically inequivalent twists

How many homotopically inequivalent $MString \wedge BU(1)^{-nT}$ are there?

Solutions: characterize homotopically inequivalent twists

How many homotopically inequivalent $MString \wedge BU(1)^{-nT}$ are there?

Theorem (Basile–Krulewski–Leone–P.-T.)

The homotopy class of $MString \wedge BU(1)^{-nT}$ only depends on the value of $n \pmod{12}$.

Solutions

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups. Smith homomorphisms provide limited success for some $-nT$ twists but we want something that works for all $-nT$ twists.
- Generators?
 - $K3 \times S^3$, $\mathbb{CP}^2 \times S^3$ are not twisted string,
 - odd dimension means no hope for complex algebraic representatives.
 - Lens spaces? Can determine twisted string structure, but hard to detect if generators.

Solutions: bordism invariants

Theorem (Basile–Krulewski–Leone–P.-T.)

Given a $(BU(1), -nT)$ -twisted string 7-manifold M , there exists a spin 8-manifold N with $\tilde{f} : N \rightarrow BU(1)$ and $\partial N = M$, then the map

$$\alpha_8^E : \Omega_8^{Spin}(BU(1)) \rightarrow \mathbb{Z}$$

$$N \mapsto \int_N \hat{A}(TN) ch(E - \text{rk } E)$$

descends to an invariant

$$\alpha_7^E(M) := \alpha_8^E(N) : \Omega_7^{String}(BU(1)^{-nT}) \rightarrow \mathbb{Q}/\mathbb{Z}$$

precisely when $ch_4(E) = nx^2 ch_2(E)$.

Theorem-in-progress (Basile–Krulewski–Leone–P.-T.)

All bordism invariants $\Omega_7^{String}(BU(1)^{-nT}) \rightarrow \mathbb{Q}/\mathbb{Z}$ arise in this way.

Solutions

- Adams spectral sequence: trivial at $p \geq 5$, extension problems at $p = 2, 3$.
- Few useful comparison maps, $\Omega_7^{\text{Spin}} \cong \Omega_7^{\text{String}} \cong 0$ and comparison maps changing the dimension land in high dimensional or 0 groups. Smith homomorphisms provide limited success for some $-nT$ twists but we want something that works for all $-nT$ twists.
- Generators?
 - $K3 \times S^3$, $\mathbb{CP}^2 \times S^3$ are not twisted string,
 - odd dimension means no hope for complex algebraic representatives.
 - Lens spaces? Can determine twisted string structure, but hard to detect if generators.

Solutions: manifold generators from sphere bundles

Instead of $\mathbb{CP}^2 \times S^3$, we consider “twisted products” $\mathbb{CP}^2 \tilde{\times} S^3$, namely sphere bundles $S(V)$ of rank 4 real vector bundles $V \rightarrow \mathbb{CP}^2$.

Solutions: manifold generators from sphere bundles

Instead of $\mathbb{CP}^2 \times S^3$, we consider “twisted products” $\mathbb{CP}^2 \tilde{\times} S^3$, namely sphere bundles $S(V)$ of rank 4 real vector bundles $V \rightarrow \mathbb{CP}^2$.

- They naturally come equipped with a map to $BU(1)$:
 $S(V) \rightarrow \mathbb{CP}^2 \rightarrow BU(1)$.

Solutions: manifold generators from sphere bundles

Instead of $\mathbb{CP}^2 \times S^3$, we consider “twisted products” $\mathbb{CP}^2 \widetilde{\times} S^3$, namely sphere bundles $S(V)$ of rank 4 real vector bundles $V \rightarrow \mathbb{CP}^2$.

- They naturally come equipped with a map to $BU(1)$:
 $S(V) \rightarrow \mathbb{CP}^2 \rightarrow BU(1)$.
- They naturally come equipped with a bulk manifold:
 $S(V) = \partial D(V)$.

Solutions: manifold generators from sphere bundles

Theorem (Basile–Krulewski–Leone–P.-T.)

The map

$$[\mathbb{CP}^2, BSO(4)] \xrightarrow{(p_1, \chi)} \mathbb{Z} \times \mathbb{Z}$$

is an isomorphism onto the subset satisfying $\binom{p_1}{2} \equiv \chi \pmod{2}$.

Solutions: manifold generators from sphere bundles

Theorem (Basile–Krulewski–Leone–P.-T.)

The map

$$[\mathbb{CP}^2, BSO(4)] \xrightarrow{(p_1, \chi)} \mathbb{Z} \times \mathbb{Z}$$

is an isomorphism onto the subset satisfying $\binom{p_1}{2} \equiv \chi \pmod{2}$.

Proof sketch.

- $[\mathbb{CP}^2, BU(2)] \xrightarrow[\sim]{(c_1, c_2)} \mathbb{Z} \times \mathbb{Z}$

Solutions: manifold generators from sphere bundles

Theorem (Basile–Krulewski–Leone–P.-T.)

The map

$$[\mathbb{CP}^2, BSO(4)] \xrightarrow{(p_1, \chi)} \mathbb{Z} \times \mathbb{Z}$$

is an isomorphism onto the subset satisfying $\binom{p_1}{2} \equiv \chi \pmod{2}$.

Proof sketch.

- $[\mathbb{CP}^2, BU(2)] \xrightarrow[\sim]{(c_1, c_2)} \mathbb{Z} \times \mathbb{Z}$
- Things in the image of $[\mathbb{CP}^2, BSU(2)] \rightarrow [\mathbb{CP}^2, BSO(4)]$ satisfy this property

Solutions: manifold generators from sphere bundles

Theorem (Basile–Krulewski–Leone–P.-T.)

The map

$$[\mathbb{CP}^2, BSO(4)] \xrightarrow{(p_1, \chi)} \mathbb{Z} \times \mathbb{Z}$$

is an isomorphism onto the subset satisfying $\binom{p_1}{2} \equiv \chi \pmod{2}$.

Proof sketch.

- $[\mathbb{CP}^2, BU(2)] \xrightarrow[\sim]{(c_1, c_2)} \mathbb{Z} \times \mathbb{Z}$
- Things in the image of $[\mathbb{CP}^2, BSU(2)] \rightarrow [\mathbb{CP}^2, BSO(4)]$ satisfy this property
- $[S^4, BSO(4)]$ acts on $[\mathbb{CP}^2, BSO(4)]$ and acting by TS^4 sends $(p, e) \rightsquigarrow (p, e + 2)$

Solutions: manifold generators from sphere bundles

Theorem (Basile–Kruewski–Leone–P.-T.)

The map

$$[\mathbb{CP}^2, BSO(4)] \xrightarrow{(p_1, \chi)} \mathbb{Z} \times \mathbb{Z}$$

is an isomorphism onto the subset satisfying $\binom{p_1}{2} \equiv \chi \pmod{2}$.

Proof sketch.

- $[\mathbb{CP}^2, BU(2)] \xrightarrow[\sim]{(c_1, c_2)} \mathbb{Z} \times \mathbb{Z}$
- Things in the image of $[\mathbb{CP}^2, BSU(2)] \rightarrow [\mathbb{CP}^2, BSO(4)]$ satisfy this property
- $[S^4, BSO(4)]$ acts on $[\mathbb{CP}^2, BSO(4)]$ and acting by TS^4 sends $(p, e) \rightsquigarrow (p, e + 2)$
- Use the fibration $S^4 \rightarrow BSO(4) \rightarrow BSO(5)$ to show TS^4 acts transitively on the set of bundles with a given p_1

Solutions: manifold generators from sphere bundles

Theorem (Basile–Krulewski–Leone–P.-T.)

The map

$$[\mathbb{CP}^2, BSO(4)] \xrightarrow{(p_1, \chi)} \mathbb{Z} \times \mathbb{Z}$$

is an isomorphism onto the subset satisfying $\binom{p_1}{2} \equiv \chi \pmod{2}$.

Theorem (Basile–Krulewski–Leone–P.-T.)

$S(V)$ equipped with $f : S(V) \xrightarrow{\pi} \mathbb{CP}^2 \subset BU(1)$ is nc_1^2 -twisted string exactly when $p_1(V) + 3 + 2n \equiv 0 \pmod{2\chi(V)}$ in which case

$$\alpha_7^E(S(V)) = \frac{ch_2(E)(p_1(V) + 3 + 2n)}{48e(V)}$$

An illustrative example: a
bordism invariant for $n = 1$

A useful cofiber sequence

$$MString \rightarrow MSpin \rightarrow MSpin/MString \rightarrow \Sigma MString$$

A useful cofiber sequence

$$(MString \rightarrow MSpin \rightarrow MSpin/MString \rightarrow \Sigma MString) \wedge BU(1)^{-T}$$

A useful cofiber sequence

$$(MString \rightarrow MSpin \rightarrow MSpin/MString \rightarrow \Sigma MString) \wedge BU(1)^{-T}$$

- $MString \wedge BU(1)^{-T}$: $(M, f : M \rightarrow BU(1), \phi)$, M spin and ϕ a string structure on $TM - f^*T$

A useful cofiber sequence

$$(MString \rightarrow MSpin \rightarrow MSpin/MString \rightarrow \Sigma MString) \wedge BU(1)^{-T}$$

- $MString \wedge BU(1)^{-T}$: $(M, f : M \rightarrow BU(1), \phi)$, M spin and ϕ a string structure on $TM - f^*T$
- $MSpin \wedge BU(1)^{-T}$: $(N, g : N \rightarrow BU(1), \psi)$, N oriented and ψ a spin structure on $TN - f^*nT$

A useful cofiber sequence

$$(MString \rightarrow MSpin \rightarrow MSpin/MString \rightarrow \Sigma MString) \wedge BU(1)^{-T}$$

- $MString \wedge BU(1)^{-T}$: $(M, f : M \rightarrow BU(1), \phi)$, M spin and ϕ a string structure on $TM - f^*T$
- $MSpin \wedge BU(1)^{-T}$: spin manifolds N with a map to $BU(1)$ ¹

¹ $MSpin \wedge BU(1)^{-T} \cong MSpin \wedge BU(1)_+$

A useful cofiber sequence

$$(MString \rightarrow MSpin \rightarrow MSpin/MString \rightarrow \Sigma MString) \wedge BU(1)^{-T}$$

- $MString \wedge BU(1)^{-T}$: $(M, f : M \rightarrow BU(1), \phi)$, M spin and ϕ a string structure on $TM - f^*T$
- $MSpin \wedge BU(1)^{-T}$: spin manifolds N with a map to $BU(1)$ ¹
- $MSpin/MString \wedge BU(1)^{-T}$: $(N, M, f : N \rightarrow BU(1), \phi)$, N is spin, $M = \partial N$, and ϕ is a string structure on $TM - f|_{\partial N}^*T$

¹ $MSpin \wedge BU(1)^{-T} \cong MSpin \wedge BU(1)_+$

Overview of the construction

We wish to construct morphisms

$$\alpha_8 : \pi_8 MSpin \wedge BU(1)^{-T} \rightarrow \mathbb{Z}$$

that descend along the diagram

$$\begin{array}{ccccc}
 \pi_8 MSpin \wedge BU(1)_+ & \rightarrow & \pi_8 MSpin / MString \wedge BU(1)^{-T} & \rightarrow & \pi_7 MString \wedge BU(1)^{-T} \\
 \downarrow & & \downarrow & & \downarrow \\
 \mathbb{Z} & \hookrightarrow & \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z}
 \end{array}$$

8d integer invariant

For $(N, f) \in \pi_8 MSpin \wedge BU(1)_+$, consider the “index of the twisted Dirac operator”

$$\alpha_8^{\mathcal{O}(1)}(N, f) = \hat{A}(TN) \text{ch}(f^* \mathcal{O}(1) - 1_{\mathbb{C}})[N]$$

8d integer invariant

For $(N, f) \in \pi_8 MSpin \wedge BU(1)_+$, consider the “index of the twisted Dirac operator”

$$\alpha_8^{\mathcal{O}(1)}(N, f) = \hat{A}(TN) \underbrace{\text{ch}(f^* \mathcal{O}(1) - 1_{\mathbb{C}})}_{\text{could twist by any } E} [N]$$

8d integer invariant

For $(N, f) \in \pi_8 MSpin \wedge BU(1)_+$, consider the “index of the twisted Dirac operator”

$$\alpha_8^{\mathcal{O}(1)}(N, f) = \hat{A}(TN) \text{ch}(f^* \mathcal{O}(1) - 1_{\mathbb{C}})[N]$$

By APS, this is a K -theory pushforward²

$$\begin{array}{ccccc} [\mathcal{O}(1) - 1_{\mathbb{C}}] \in KU(BU(1)) & \xrightarrow{f^*} & KU(N) & \xrightarrow{i!} & KU(\text{pt}) \cong \mathbb{Z} \\ & & \hat{A}(N) \text{ch} \downarrow & & \downarrow \\ & & H(N; \mathbb{Q}) & \xrightarrow{\int_N} & H(\text{pt}; \mathbb{Q}) \cong \mathbb{Q} \end{array}$$

²Can also define as a KO pushforward

8d integer invariant

For $(N, f) \in \pi_8 MSpin \wedge BU(1)_+$, consider the “index of the twisted Dirac operator”

$$\alpha_8^{\mathcal{O}(1)}(N, f) = \hat{A}(TN) \text{ch}(f^* \mathcal{O}(1) - 1_{\mathbb{C}})[N]$$

By APS, this is a K -theory pushforward²

$$\begin{array}{ccccc} [\mathcal{O}(1) - 1_{\mathbb{C}}] \in KU(BU(1)) & \xrightarrow{f^*} & KU(N) & \xrightarrow{i_!} & KU(\text{pt}) \cong \mathbb{Z} \\ & & \hat{A}(N) \text{ch} \downarrow & & \downarrow \\ & & H(N; \mathbb{Q}) & \xrightarrow{\int_N} & H(\text{pt}; \mathbb{Q}) \cong \mathbb{Q} \end{array}$$

So it defines an integer-valued map

$$\alpha_8^{\mathcal{O}(1)} : \pi_8 MSpin \wedge BU(1)^{-nT} \rightarrow \mathbb{Z}$$

²Can also define as a KO pushforward

8d integer invariant

Claim

$\alpha_8^{\mathcal{O}(1)}$ can take on any integer value.

8d integer invariant

Claim

$\alpha_8^{\mathcal{O}(1)}$ can take on any integer value.

Proof. Consider a degree 4 hypersurface $N \subset \mathbb{CP}^5$, it naturally comes equipped with $f : N \subset \mathbb{CP}^5 \rightarrow \mathbb{CP}^\infty$.

8d integer invariant

Claim

$\alpha_8^{\mathcal{O}(1)}$ can take on any integer value.

Proof. Consider a degree 4 hypersurface $N \subset \mathbb{CP}^5$, it naturally comes equipped with $f : N \subset \mathbb{CP}^5 \rightarrow \mathbb{CP}^\infty$. We have $TN + \mathcal{O}(1) = 6\mathcal{O}(1) - \mathcal{O}(4)$, so

$$p_1(TN) = (6 - 16)x^2 = -10x^2$$

8d integer invariant

Claim

$\alpha_8^{\mathcal{O}(1)}$ can take on any integer value.

Proof. Consider a degree 4 hypersurface $N \subset \mathbb{CP}^5$, it naturally comes equipped with $f : N \subset \mathbb{CP}^5 \rightarrow \mathbb{CP}^\infty$. We have $TN + \mathcal{O}(1) = 6\mathcal{O}(1) - \mathcal{O}(4)$, so

$$p_1(TN) = (6 - 16)x^2 = -10x^2$$

It follows that

$$\begin{aligned}\alpha_8^{\mathcal{O}(1)}(N, f) &= \int_N 10x^4/48 + x^4/24 = \int_N 12x^4/48 \\ &= \int_{\mathbb{P}^5} 4(12h^4/48) = 1\end{aligned}$$

8d integer invariant

The invariant can be written as

$$\begin{aligned}\alpha_8^{\mathcal{O}(1)}(N, f) &= \hat{A}(TN) \text{ch}(f^* \mathcal{O}(1) - 1_{\mathbb{C}})[N] \\ &= - \int_N (p_1(TN) - 2x^2)x^2/48 \\ &= - \int_N p_1(TN - f^* T)x^2/48\end{aligned}$$

Extending to the relative groups

Take $(N, M, f, \phi) \in \pi_8 MSpin / MString \wedge BU(1)^{-T}$,

$$\begin{array}{ccc}
 & & BString \\
 & \nearrow \phi & \downarrow \\
 M & \xrightarrow{TM-f|_{\partial N}^* T} & BSpin
 \end{array}$$

Extending to the relative groups

Take $(N, M, f, \phi) \in \pi_8 MSpin / MString \wedge BU(1)^{-T}$,

$$\begin{array}{ccc}
 & & BString \\
 & \nearrow \phi & \downarrow \\
 M & \xrightarrow{TM - f|_{\partial N}^* T} & BSpin
 \end{array}$$

- $BString$ is 7-connected, so we may choose a trivialization $\tilde{\phi}$ of $TM - f|_{\partial N}^* T$

Extending to the relative groups

Take $(N, M, f, \phi) \in \pi_8 MSpin / MString \wedge BU(1)^{-T}$,

$$\begin{array}{ccc}
 & & BString \\
 & \nearrow \phi & \downarrow \\
 M & \xrightarrow{TM - f|_{\partial N}^* T} & BSpin
 \end{array}$$

- $BString$ is 7-connected, so we may choose a trivialization $\tilde{\phi}$ of $TM - f|_{\partial N}^* T$
- This defines a relative KO class $[TN - f^* T]_{\tilde{\phi}} \in KO(N, M)$ and a relative class $p_1([TN - f^* T]_{\tilde{\phi}}) \in H^4(N, M)$

Extending to the relative groups

Take $(N, M, f, \phi) \in \pi_8 MSpin / MString \wedge BU(1)^{-T}$,

$$\begin{array}{ccc}
 & & BString \\
 & \nearrow \phi & \downarrow \\
 M & \xrightarrow{TM - f|_{\partial N}^* T} & BSpin
 \end{array}$$

- $BString$ is 7-connected, so we may choose a trivialization $\tilde{\phi}$ of $TM - f|_{\partial N}^* T$
- This defines a relative KO class $[TN - f^* T]_{\tilde{\phi}} \in KO(N, M)$ and a relative class $p_1([TN - f^* T]_{\tilde{\phi}}) \in H^4(N, M)$
- Define $\alpha_{\text{rel}}^{\mathcal{O}(1)} = -p_1([TN - f^* T]_{\tilde{\phi}})c_1(\mathcal{L})^2[N]$ using Poincare-Lefschetz duality

Extending to the relative groups

$$\begin{array}{ccc} \pi_8 MSpin \wedge BU(1)_+ & \longrightarrow & \pi_8 MSpin / MString \wedge BU(1)^{-T} \\ \downarrow \alpha_8^{\mathcal{O}(1)} & & \downarrow \alpha_{\text{rel}}^{\mathcal{O}(1)} \\ \mathbb{Z} & \hookrightarrow & \mathbb{Q} \end{array}$$

Descending to a 7d invariant

$$\begin{array}{ccc} \pi_8 MSpin/MString \wedge BU(1)^{-T} & \longrightarrow & \pi_7 MString \wedge BU(1)^{-T} \\ \downarrow \alpha_{\text{rel}}^{\mathcal{O}(1)} & & \downarrow \alpha_7^{\mathcal{O}(1)} \\ \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z} \end{array}$$

Descending to a 7d invariant

$$\begin{array}{ccc} \pi_8 MSpin / MString \wedge BU(1)^{-T} & \longrightarrow & \pi_7 MString \wedge BU(1)^{-T} \\ \downarrow \alpha_{\text{rel}}^{\mathcal{O}(1)} & & \downarrow \alpha_7^{\mathcal{O}(1)} \\ \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z} \end{array}$$

- $\phi : M \xrightarrow{TM - f^*T} BString$ lifts to a framing of $TM - f^*T$

Descending to a 7d invariant

$$\begin{array}{ccc}
 \pi_8 MSpin / MString \wedge BU(1)^{-T} & \twoheadrightarrow & \pi_7 MString \wedge BU(1)^{-T} \\
 \downarrow \alpha_{\text{rel}}^{\mathcal{O}(1)} & & \downarrow \alpha_7^{\mathcal{O}(1)} \\
 \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z}
 \end{array}$$

- $\phi : M \xrightarrow{TM - f^*T} BString$ lifts to a framing of $TM - f^*T$
- $\pi_7(MSpin \wedge BU(1)_+) \simeq 0$ so there exists a pair $(N^8, \tilde{f})$ with $(M, f) = \partial(N, \tilde{f})$

Descending to a 7d invariant

$$\begin{array}{ccc}
 \pi_8 MSpin / MString \wedge BU(1)^{-T} & \twoheadrightarrow & \pi_7 MString \wedge BU(1)^{-T} \\
 \downarrow \alpha_{\text{rel}}^{\mathcal{O}(1)} & & \downarrow \alpha_7^{\mathcal{O}(1)} \\
 \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z}
 \end{array}$$

- $\phi : M \xrightarrow{TM - f^*T} BString$ lifts to a framing of $TM - f^*T$
- $\pi_7(MSpin \wedge BU(1)_+) \simeq 0$ so there exists a pair $(N^8, \tilde{f})$ with $(M, f) = \partial(N, \tilde{f})$
- Set $\alpha_7^{\mathcal{O}(1)}(M, f, \phi) := \alpha_{\text{rel}}^{\mathcal{O}(1)}(N, M)$

Summary

$$\begin{array}{ccccc} \pi_8 MSpin \wedge BU(1)_+ & \longrightarrow & \pi_8 MSpin / MString \wedge BU(1)^{-T} & \longrightarrow & \pi_7 MString \wedge BU(1)^{-T} \\ \downarrow & & \downarrow & & \downarrow \\ \mathbb{Z} & \hookrightarrow & \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z} \end{array}$$

Summary

$$\begin{array}{ccccc}
 \pi_8 MSpin \wedge BU(1)_+ & \longrightarrow & \pi_8 MSpin / MString \wedge BU(1)^{-T} & \longrightarrow & \pi_7 MString \wedge BU(1)^{-T} \\
 \downarrow & & \downarrow & & \downarrow \\
 \mathbb{Z} & \hookrightarrow & \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z}
 \end{array}$$

Index of Dirac operator twisted by $\mathcal{O}(1) - 1_{\mathbb{C}}$ is a K -theory pushforward.

Summary

$$\begin{array}{ccccc}
 \pi_8 MSpin \wedge BU(1)_+ & \rightarrow & \pi_8 MSpin / MString \wedge BU(1)^{-T} & \rightarrow & \pi_7 MString \wedge BU(1)^{-T} \\
 \downarrow & & \downarrow & & \downarrow \\
 \mathbb{Z} & \hookrightarrow & \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z}
 \end{array}$$

It factors a $p_1([TN - f^*T])$ so it extends to a relative invariant.

Summary

$$\begin{array}{ccccc}
 \pi_8 MSpin \wedge BU(1)_+ & \rightarrow & \pi_8 MSpin / MString \wedge BU(1)^{-T} & \rightarrow & \pi_7 MString \wedge BU(1)^{-T} \\
 \downarrow & & \downarrow & & \downarrow \\
 \mathbb{Z} & \hookrightarrow & \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z}
 \end{array}$$

Any twisted string 7-manifold is on the boundary of a spin 8-manifold so this descends to a 7d invariant.

Summary

$$\begin{array}{ccccc} \pi_8 MSpin \wedge BU(1)_+ & \longrightarrow & \pi_8 MSpin / MString \wedge BU(1)^{-nT} & \longrightarrow & \pi_7 MString \wedge BU(1)^{-nT} \\ \downarrow & & \downarrow & & \downarrow \\ \mathbb{Z} & \hookrightarrow & \mathbb{Q} & \longrightarrow & \mathbb{Q}/\mathbb{Z} \end{array}$$

General case: need to choose E to twist Dirac operator by that both descends and detects the most torsion.

Manifold generators

Twisted string structures on $S(V)$

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with $p = p_1(V_{p,\chi})$,
 $\chi = e(V_{p,\chi})$, $\binom{p}{2} \equiv \chi \pmod{2}$

Twisted string structures on $S(V)$

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with $p = p_1(V_{p,\chi})$,
 $\chi = e(V_{p,\chi})$, $\binom{p}{2} \equiv \chi \pmod{2}$

- When is $D(V_{p,\chi})$ spin?

Twisted string structures on $S(V)$

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with $p = p_1(V_{p,\chi})$,
 $\chi = e(V_{p,\chi})$, $\binom{p}{2} \equiv \chi \pmod{2}$

- When is $D(V_{p,\chi})$ spin? Exactly when p is odd, in which case it has a unique spin structure.

Twisted string structures on $S(V)$

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with $p = p_1(V_{p,\chi})$,
 $\chi = e(V_{p,\chi})$, $\binom{p}{2} \equiv \chi \pmod{2}$

- When is $D(V_{p,\chi})$ spin? Exactly when p is odd, in which case it has a unique spin structure.
- When does $S(V_{p,\chi})$ admit a $(BU(1), -T)$ -twisted string structure?

Twisted string structures on $S(V)$

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with $p = p_1(V_{p,\chi})$,
 $\chi = e(V_{p,\chi})$, $\binom{p}{2} \equiv \chi \pmod{2}$

- When is $D(V_{p,\chi})$ spin? Exactly when p is odd, in which case it has a unique spin structure.
- When does $S(V_{p,\chi})$ admit a $(BU(1), -T)$ -twisted string structure?

$$KO(D(V), S(V)) \rightarrow \widetilde{KO}(D(V)) \ni [TD(V) - f^*T]$$

Twisted string structures on $S(V)$

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with $p = p_1(V_{p,\chi})$,
 $\chi = e(V_{p,\chi})$, $\binom{p}{2} \equiv \chi \pmod{2}$

- When is $D(V_{p,\chi})$ spin? Exactly when p is odd, in which case it has a unique spin structure.
- When does $S(V_{p,\chi})$ admit a $(BU(1), -T)$ -twisted string structure?

$$KO(D(V), S(V)) \rightarrow \widetilde{KO}(D(V)) \ni [TD(V) - f^*T]$$

$$[TD(V) - f^*T] = [V + 3\mathcal{O}(1) - (\mathcal{O}(1) + \mathcal{O}(-1))]$$

$\alpha_7^{\mathcal{O}(1)}$ on $-T$ -twisted sphere bundles

Theorem (Basile–Krulewski–Leone–P.-T.)

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with
 $p = p_1(V_{p,\chi})$, $\chi = e(V_{p,\chi})$, $S(V_{p,\chi})$ admits a $(BU(1), -T)$ -twisted
 string structure exactly when

$$p + 1 \equiv 0 \pmod{2\chi}$$

in which case

$$\alpha_7(S(V_{p,\chi})) = -\frac{p+1}{48\chi}$$

α_7^E on $-nT$ -twisted sphere bundles

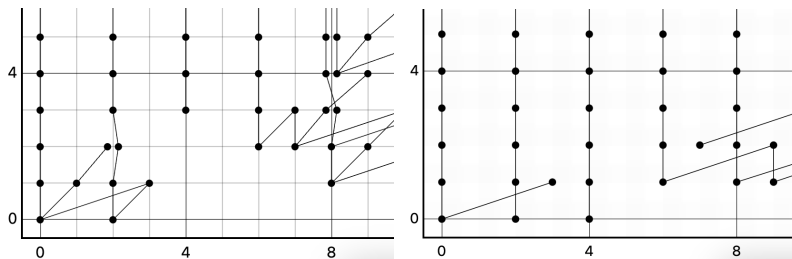
Theorem (Basile–Krulewski–Leone–P.-T.)

Given a rank 4 bundle $V_{p,\chi} \xrightarrow{p} \mathbb{CP}^2$ with $p = p_1(V_{p,\chi})$, $\chi = e(V_{p,\chi})$, $S(V_{p,\chi})$ admits a $(BU(1), -T)$ -twisted string structure exactly when

$$p + 3 - 2n \equiv 0 \pmod{2\chi}$$

in which case

$$\alpha_7^E = -\frac{ch_2(E)(p + 3 - 2n)}{48\chi}$$

Order of the $n = 1$ bordism group

- $p = 3, \chi = 1$ satisfies $\binom{3}{2} \equiv 1 \pmod{2}$, $3 + 1 \equiv 0 \pmod{2}$ and

$$\alpha_7^{\mathcal{O}(1)}(S(V_{3,1})) = -\frac{1}{12}$$

Thank you!